# Supersymmetric candidates for Dark Matter

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Steffen, arXiv:0811.3347 Hooper, arXiv:0901.4090 March 24<sup>th</sup> 2009

ICC - Journal Club

The evidence for Dark Matter

- There are many hints at different scales
  - galaxy rotation curves





The evidence for Dark Matter

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  - cluster of galaxies



#### Clusters of galaxies

- or rotation curves
- gravitational lensing
   Light bends differently than
   predicted from GR, if only luminous
   matter is taken into account

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#### Dark Matter exists!

It consists of particles Permeates galactic haloes

So, what is the Dark Matter?

The SM do not provide a good candidate for Dark Matter.

SM neutrinos could be a candidate but they are too light (HDM) implying top-down scenarios not supported by the present observations (the galaxies seems older than clusters).

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There are other candidates beyond the Standard model?

- Heavy neutrinos Lee & Weinberg, 77; Dolgov, 02; Griest & Kamionkowski, 90
- Sterile neutrinos Dodelson & Widrow, 93; Abazajian, Fuller & Patel, 01
- Axions Peccei & Quinn, 97; Rosenberg & van Bibber, 00
- ► SUSY (LSP)
- Extra dimensions (LKP) Cheng, Feng & Matchev, 02; Agashe & Servant, 04
- Scalar Dark Matter Boehm, Fayet & Silk, 04
- Dark Matter form little Higgs model Birkedal & Wacker, 03



The MSSM is the minimal supersymmetrical extension of the SM Additionally to solve

- Hierarchy problem
- Unification of gauge couplings

The MSSM with R-parity could provide a good candidate for DM.



#### **R**-parity

R-parity introduced to avoid

- baryon number and lepton number violation
- too fast proton decay

$$R = (-1)^{3B + L + 2S}$$

- R = +1 for ordinary particles
- R = -1 for SUSY particles

Important phenomenological consequences

- SUSY particles always produced in pairs
- a SUSY particle decays in an odd number of SUSY particles
- the lightest SUSY particle is stable

Quarks	$u_{R,L}$ $d_{R,L}$
	C <sub>R,L</sub> S <sub>R,L</sub>
	$t_{R,L}$ $b_{R,L}$
Leptons	$e_{R,L}$ $v_e$
	$\mu_{R,L}$ $\nu_{\mu}$
	$ au_{R,L}$ $ u_{ au}$
Gauge bosons	$Z^0, W^{\pm}, g$
Squarks	$\tilde{u}_{R,L}$ $\tilde{d}_{R,L}$
	$ ilde{c}_{R,L}$ $ ilde{s}_{R,L}$
	$\tilde{t}_{R,L}$ $\tilde{b}_{R,L}$
Sleptons	$\tilde{e}_{R,L}$ $\tilde{v}_{e}$
	$ ilde{\mu}_{R,L}$ $ ilde{ u}_{\mu}$
	$ ilde{ au}_{R,L}$ $ ilde{ u}_{ au}$
Charginos	$\chi_1^{\pm}, \chi_2^{\pm}$
Neutralinos	$\chi_1^0, \chi_2^{\overline{0}}$
	$\chi_{3}^{0}, \chi_{4}^{\overline{0}}$
Gluinos	ĝ

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- In SUGRA models
  - gravitino
- MSSM + PQ mechanism
  - 🕨 axino

### Neutralinos in the MSSM

Neutralinos in the MSSM are physical superpositions of the gauginos (bino  $\tilde{B}$  and wino  $\tilde{W}$ ) and the Higgsinos ( $\tilde{H}_u$  and  $\tilde{H}_d$ )

Neutralinos are WIMPs, thermally produced.

Mass matrix

$$\mathcal{M}_{\chi^{0}} = \begin{pmatrix} M_{1} & 0 & -M_{Z} \, s_{\theta} \, c_{\beta} & M_{Z} \, s_{\theta} \, s_{\beta} \\ 0 & M_{2} & M_{Z} \, c_{\theta} \, c_{\beta} & M_{Z} \, c_{\theta} \, s_{\beta} \\ -M_{Z} \, s_{\theta} \, c_{\beta} & M_{Z} \, c_{\theta} \, c_{\beta} & 0 & -\mu \\ -M_{Z} \, s_{\theta} \, s_{\beta} & M_{Z} \, c_{\theta} \, s_{\beta} & -\mu & 0 \end{pmatrix}$$

 $M_i$ : gaugino masses, break SUSY softly  $\mu$ : higgsino mass, appear in the superpotential

#### Lightest neutralino

$$\chi_1^0 = N_{11} \, \tilde{B}^0 + N_{12} \, \tilde{W}^0 + N_{13} \, \tilde{H}_d^0 + N_{14} \, \tilde{H}_u^0$$

The detection properties of the lightest neutralino depend on its composition

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$$\begin{array}{ccc} M_Z \, s_\theta \, c_\beta & M_Z \, s_\theta \, s_\beta \\ M_Z \, c_\theta \, c_\beta & M_Z \, c_\theta \, s_\beta \\ 0 & -\mu \\ -\mu & 0 \end{array} \right)$$

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Let us recall that  $M_1$ ,  $M_2$  and  $\mu$  are not (usually) free parameters

•  $\mu$  fixed by the requirement of a proper radiative EW symmetry breaking

$$\mu^{2}(M_{Z}) = \frac{m_{H_{2}}^{2} \sin^{2}\beta - m_{H_{1}}^{2} \cos^{2}\beta}{\cos 2\beta} - \frac{M_{Z}^{2}}{2}$$

•  $M_1$  and  $M_2$  fixed at a grand unification scale ( $M_{GUT} \sim 2 \cdot 10^{26}$  GeV) and then evolved down to the EW scale by the use of the RGE

$$\rightarrow \quad M_1 = \frac{5}{3} \tan^2 \theta_W \sim M_2/2 \text{ at } Q = M_Z$$



The most general MSSM contains  $\sim$  120 free parameters. We need a phenomenologically viable model!

#### Constraining the MSSM

The constraint MSSM contains 5 free parameters

- $\tan\beta$ : ratio of the vevs
- sign(μ)
- m<sub>1/2</sub>: Universal gaugino mass
- m<sub>0</sub>: Universal scalar (sfermions and Higgs bosons) mass
- A<sub>0</sub>: Universal trilinear couplings



Steffen, 08; Hooper, 09

#### Bulk region

- \* Neutralino is bino-like
- \* Light sleptons
- $\rightarrow$  annihilation via slepton exchange into a lepton pair:

$$\chi_1^0\chi_1^0\to I^+\,I^-$$

> But annihilation cross section small and then  $\Omega h^2$  too large



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 But annihilation cross section small and then Ωh<sup>2</sup> too large

We need to enlarge the annihilation cross section!



Steffen, 08; Hooper, 09

#### Focus point

\* Neutralino is a Bino-Higgsino mixture

\* Annihilation into vector bosons becomes efficient

$$\chi_1^0 \chi_1^0 \rightarrow W^+ W^-, Z^0 Z^0$$

Vector bosons only couple to Higgsino-like neutralinos!

Good relic density



#### Steffen, 08; Hooper, 09

#### **Coannihilation region**

\* If LSP and NLSP almost degenerate in mass \* Important enhancement for the cross-section

$\chi_1^0  \tilde{\tau} \to \tau  h,  \tau  Z,  \tau  \gamma$	$\sigma \propto Exp[-\Delta_m]$
$\tilde{\tau}\tilde{\tau} \rightarrow \tau\tau,hh,\gamma\gamma$	$\sigma \propto Exp[-2\Delta_m]$

 $\Delta_m \equiv (m_{NLSP} - m_{LSP})/m_{LSP}$ Griest & Seckel, 91

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Coannihilation could also take place with any other SUSY particle:  $\chi^{\pm}, \chi^{0}, \tilde{g}...$ 



Steffen, 08; Hooper, 09

**Higgs funnel** (for large  $tan \beta$ )

$$M_A^2 \propto -\frac{1}{\cos 2\beta} B \mu$$

\* In the region when  $M_A \sim 2 M_{\chi_1^0}$ \* Resonance in the *s*-channel

$$\chi_1^0\chi_1^0\to A\to b\,\bar{b}$$

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Note that it is also possible to have resonances with h or Z in the general MSSM

### Neutralinos: Direct detection

#### Spin-independent cross section



#### Higgs exchange

$$\sigma_{\chi_1^0 - p} \propto rac{m_r^4}{4\pi} \, rac{y_q^2}{m_h^4} \left| N_{13, \, 14} \, \left( g' \, N_{11} - g \, N_{12} 
ight) 
ight|^2$$

It is the leading contribution. Maximum enhancement for:

- Higgsino-Íike neutralino  $(\mu \ll M_1)$
- Iow Higgs mass



#### squark exchange

$$\sigma_{\chi_1^0 - \rho} \propto \frac{m_r^2}{4\pi} \left( \frac{g'^2 \sin \theta_W}{m_{\tilde{q}}^2 - m_\chi^2} \right)^2 \, |N_{11}|^4$$

• only relevant for a bino-like neutralino  $(M_1 \ll \mu)$ 

### Neutralinos: Direct detection

#### Spin-independent cross section



Exclusion regions for different SUSY models

- Baltz & Gondolo, 03
- Baltz & Gondolo, 04
- Ellis, Feng, Ferstl, Matchev & Olive, 02
- Battaglia et al., 04
- Roszkowski, Ruiz de Austri & Trotta, 07

### Sneutrinos in the MSSM

In the MSSM sneutrinos are pure left-handed! Pure left-handed sneutrinos are WIMPs thermally produced. They faces some problems...



Sneutrino (left-handed) coupling with *Z* boson is rather large, leading to

➤ Too large annihilation cross-section → implying too small relic density

Ibáñez, 84; Ellis, Hagelin, Nanopoulos & Olive, 84; Hagelin, Kane & Rabi, 84; Goodmann & Witten, 85; Freese, 86

 Too large direct detection cross section already disfavoured by current experiments Falk, Olive & Srednicki, 94

## Sneutrinos in the MSSM



Too large annihilation cross-section Arina & Fornengo, 07; Arina, 08 Too large direct detection cross section

> These problems can be alleviated by reducing the  $Z \tilde{v} \tilde{v}$  coupling

This can be done by including a 'sterile' (e.g. right-handed sneutrino) component with which the left-handed sneutrino can mix!!!

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- smaller annihilation cross section
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However in SUGRA sneutrino mixing is proportional to small neutrino Yukawas  $\rightarrow$  No enough mixing



Model independent:

$$\tilde{\nu} = \cos\theta \,\tilde{\nu}_L + \sin\theta \,\tilde{N}$$

letting  $\theta$  as a free parameter

Arina & Fornengo, 07; Arina, 08

### Pure right-handed Sneutrinos

Alternatively, a pure right-handed sneutrino could give rise to the good relic density (no coupling with Z boson)



 However: in general very small detection cross section (would not account for a WIMP observation)

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Right-handed sneutrinos could explain **PAMELA** and **ATIC** data! Allahverdi, Dutta, Richardson-McDaniel & Santoso, 09

## Gravitino

- \* Assuming that SUSY is realized not only as global but also as a local symmetry:  $\tilde{G}$  with spin 3/2
- \* Gravitino appears in gauge-mediated and gravity-mediated SUSY breaking schemes
- \* Gravitino is a singlet with respect to the gauge groups of the SM

Gravitino is a very weakly interacting particle Interactions typically given by dim 5 operators  $\rightarrow$  suppressed by  $1/M_P$ 

- Practically undetectable (except for this gravitational effect)
- > NLSP could be long-lived!  $\tau_{NLSP} \sim 1 \text{ s}$

Many possibilities for the NLSP ( $\tilde{\chi}_1^0, \tilde{\nu}, \tilde{t}...$ ) each with its own phenomenology

### Gravitino

Gravitino DM relic density

$$\Omega_{ ilde{G}}h^2 = \Omega_{ ilde{G}}^T h^2 + \Omega_{ ilde{G}}^{NT} h^2$$

High decoupling temperature because of they extremely weak interactions.

Thermal: 
$$\Omega_{\tilde{G}}^{T}h^{2} \simeq 0.27 \left(\frac{T_{R}}{10^{10} \text{ GeV}}\right) \left(\frac{100 \text{ GeV}}{m_{\tilde{G}}}\right) \left(\frac{m_{1/2}}{1 \text{ TeV}}\right)^{2}$$

Non-thermal production comes from the decay NLSP  $\rightarrow$  LSP + X

Non-thermal: 
$$\Omega_{\tilde{G}}^{NT}h^2 = \frac{m_{3/2}}{m_{NLSP}}\Omega_{NLSP}h^2$$

## Gravitino in the cMSSM



Regions with the correct relic density:

- \* only non-thermal production vs.
- \* thermal and non-thermal production

Non-thermal production alone is not sufficient. Large contributions from thermal processes are necessary.

## Conclusions

- The existence of dark matter provides strong evidence for physics beyond the Standard Model.
- The lightest neutralino is the most promising.
   Could be tested by direct and indirect detection experiments but also in colliders as pure missing *E*<sub>T</sub>.
- Pure left-handed sneutrino is already excluded, but a mixing with a right-handed sneutrino could give rise to a good DM candidate.
- For gravitinos, Dark matter experiments would not detect anything! However, in accelerators (LHC): detection of a (meta)-stable and electrically charged NLSP (τ̃).

## **Relic Density**

Principle of 'Thermal Relics' (Simplest scenario)



- Begin from a state of thermodynamical equilibrium.
- The universe expands → its temperature falls.
   SM particles no longer energetic enough to produce *χχ* pairs.
- >  $\chi$ 's density falls asymptotically.

Evolution described by the Boltzmann equation

$$\frac{dn}{dt} + 3 H n = -\langle \sigma v \rangle (n^2 - n_{eq}^2)$$

A gross estimation of a particle's relic density can be provided by

$$\Omega h^2 \sim \frac{3 \cdot 10^{-27} \, cm^3 \, s^{-1}}{\langle \sigma v \rangle} \sim \frac{1}{\langle \sigma v \rangle} \sim \frac{m_{\chi}^2}{g_{\chi}^4}$$

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The most general MSSM contains  $\sim$  120 free parameters. We need a phenomenologically viable model!

#### Constraining the MSSM

- All the soft SUSY-breaking parameters are real
- The matrices for the sfermions masses are diagonal
- The matrices for the trilinear couplings are diagonal
- The soft sfermion masses and trilinear couplings are equal for the 1<sup>st</sup> and 2<sup>nd</sup> generation
- Universality at M<sub>GUT</sub>



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## **Direct Detection**

The direct detection of Dark Matter can take place through their interaction with nuclei inside a detector



#### Problems

The nuclear recoiling energy is measured

- Ionization on solids
- Ionization in scintillators (measured by the emitted photons)
- Temperature increase (measured by the released phonons)

- Very small interaction rate
- Large backgrounds (experiments must be deep underground)
- Uncertainties in the DM properties in our galaxy

## Neutralinos and PAMELA

Two main options:

A 200 GeV wino-like neutralino annihilating into  $W^+ W^-$  could explain Pamela data

- But: ★ Wino-like neutralinos usually coannihilates with the lightest chargino generating not enough relic density
  - X A large variation on the rate of energy loss of the positrons is needed
  - ✗ Not possible to explain at the same time Pamela and Atic data

Grajek, Kane, Phalen, Pierce & Watson, 08

A 600 – 1000 GeV wino-like neutralino annihilating into  $W^+ W^-$ 

But: ★ a large nearby (~ 1 – 2 kiloparsecs of the Solar System) clump of annihilating neutralinos is needed

Hooper, Stebbins & Zurek, 08

## Gravitino Dark Matter

